





SmEdA Statistical Modal Energy distribution Analysis

Laurent MAXIT Laboratoire Vibrations-Acoustique - INSA Lyon laurent.maxit@insa-lyon.fr Vibro-acoustic modelling of complex mechanical structures in the mid-frequency range



LVA colleagues involved in SmEdA: Jean-Louis Guyader, Nicolas Totaro, Kerem Ege, HaDong Hwang, Youssef Gerges.

L. Maxit – Reformulation and extension of SEA model by taking the modal energy distribution into account, PhD thesis, INSA-Lyon, 2000 (in French!).

Outline of the presentation

- I Dual Modal Formulation
- II Fundamentals of SmEdA
- III Interests of SmEdA
- IV Extension to non-resonant transmission
- V Methodology for including dissipative threatments
- VI Modelling of the vibration transmission through industrial structures



Two coupled subsystems excited by white noise forces in $[\Omega_1, \Omega_2]$

Modal coupling schema suggested by SEA

$$\begin{cases} M_{p} \Big[a_{p}''(t) + \Delta_{p} a_{p}'(t) + \omega_{p}^{2} a_{p}(t) \Big] - \sum_{r=1}^{N_{2}} W_{pr} b_{r}'(t) = F_{p}(t), \quad \forall p \in [1, N_{1}] \\ M_{q} \Big[b_{p}''(t) + \Delta_{p} b_{p}'(t) + \omega_{p}^{2} b_{p}(t) \Big] + \sum_{s=1}^{N_{1}} W_{sq} a_{s}'(t) = F_{q}(t), \quad \forall q \in [1, N_{2}] \end{cases}$$

Subsystem modes
Oscillators
$$Mode \text{ couplings with the others subsystem modes} \qquad \text{Resonant modes}$$



Two coupled subsystems excited by white noise forces in $[\Omega_1, \Omega_2]$

Modal coupling schema suggested by SEA

$$\begin{cases} M_1 \Big[a_1''(t) + \Delta_1 a_1'(t) + \omega_1^2 a_1(t) \Big] & -G_c b_2'(t) = F_1(t) \\ M_2 \Big[b_2''(t) + \Delta_2 b_2'(t) + \omega_2^2 b_2(t) \Big] + G_c a_1'(t) = F_2(t) \end{cases}$$

$$P_{12} = \beta \left(E_1 - E_2\right)$$

→ SEA or SmEdA models



[*] L. Maxit, J.L. Guyader - Estimation of sea coupling loss factors using a dual formulation and fem modal information, part 1 : theory. *Journal of Sound and Vibration,* 239 (2001) 907-930.





Second example



4 panels coupled together at right angle excited by a point force.

Case	Floor material	Floor thickness	Wall material	Wall thickness
1	Concrete	0.2 m	Brickwork	0.04 m
2	Concrete	0.2 m	Concrete	0.2 m

SEA weak coupling assumption (Fahy and James, JSV 190 (1996)):

"Under the conditions of weak coupling, the system modes are 'localised' in the sense that **they closely resemble in natural frequency and shape the modes of the uncoupled subsystems** (given the appropriate boundary conditions)...

... Depending on the nature of the coupling, the boundary conditions for the uncoupled system do not always correspond to free displacement at the coupling".



Illustration of the numerical process for intermodal work estimation:





Intermodal works for resonant modes in the 500 Hz Octave band

Subsystem energy response in function of the frequency. Test case n°1.



Comparison between reference (full) and DMF with Non-Resonant (NR) modes (dash)

Comparison between DMF with NR modes (dash) and DMF without NR modes (dotted)

		1+2	3	4
	Panel			
	Energy			
Test case 1	Reference	79.0 dB	58.9 dB	55.0 dB
		(78.2+71.6)		
	DMF	78.3 dB	58.2 dB	54.8 dB
	DMF (resonant modes only)	78.3 dB	57.9 dB	54.2 dB
Test case 2	Reference	78.6 dB	67.2 dB	68.0 dB
		(77.8+70.9)		
	DMF	78.8 dB	45.2 dB	45.6 dB
	DMF (resonant modes only)	78.8 dB	42.9 dB	42.5 dB

Comparison of subsystem energies obtained with different calculations. Results for the octave band 500 Hz (dB, ref. 10⁻¹²J).

L. Maxit, Investigation for obtaining a modal coupling scheme in agreement with the SmEdA/SEA models for multisubsystems connected at a junction, *Proceedings of NOVEM 2015*, Dubrovnik, Croatia, April 2015. 22 p.

Third example: a rectangular plate coupled to a cavity

- Case of the cavity filled of air
 - → Good description with DMF taking only the resonant modes
- Case of the cavity filled of water (heavy fluid)

Comparison between reference (full), DMF without Non-Resonant modes (dotted), and DMF with Non-Resonant (NR) modes (dash)

L. Maxit, Analysis of the modal energy distribution of an excited vibrating panel coupled with a heavy fluid cavity by a dual modal formulation. *Journal of Sound and Vibration*, vol. 332, p. 6703-6724, 2013.

A mechanical **impedance mismatch** between the two subsystems is need.

In this case, the **soft** subsystem is represented by its uncoupled-**blocked** modes whereas the other one (i.e. the **stiffer** one) is represented by its uncoupled-**free** modes.

In the following, we suppose to be in these conditions.

SmEdA is based on:

- the modal description of uncoupled subsystems (natural frequencies, mode shapes)
- the same assumptions as SEA except for the modal energy equipartition
- the description of the energy sharing between modes rather than between subsystems

• Time-averaged power for a single oscillator

An oscillator excited by a white noise force

Equation of motions: $M[\ddot{x}(t) + \omega_0 \eta_0 \dot{x}(t) + (\omega_0)^2 x(t)] = F(t)$

Time-averaged energy (kinetic + strain):

$$\langle E_{t} \rangle_{t} = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{+T} \left\{ \frac{1}{2} M \left[\dot{x}(t) \right]^{2} + \frac{1}{2} K \left[x(t) \right]^{2} \right\} dt$$

(when time-averaged)

$$\langle P_{diss} \rangle_{t} = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{+T} \omega_{0} \eta_{0} [\dot{x}(t)]^{2} dt \qquad \Rightarrow \qquad \langle P_{diss} \rangle_{t} \approx \omega_{0} \eta_{0} \langle E_{t} \rangle_{t}$$
(1)
$$\langle P_{inj} \rangle_{t} = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{+T} F(t) \dot{x}(t) dt \qquad \Rightarrow \qquad \langle P_{inj} \rangle_{t} = \frac{\pi \overline{S}_{F}}{4M}$$
(2)

R.H. Lyon, R.G. DeJong *Theory and Application of Statistical Energy Analysis,* Butterworth-Heineman (1995)

 Time-averaged power exchanged by two oscillators: Case of white noise forces in [0,∞[

Two oscillators coupled by a gyrostatic element and excited by uncorrelated white noise forces Equation of motions $\begin{cases}
\ddot{x}_{1}(t) + \Delta_{1}\dot{x}_{1}(t) + \omega_{1}^{2}x_{1}(t) - \sqrt{M_{1}^{-1}M_{2}} \ \gamma \ \dot{x}_{2}(t) = F_{1}(t), \\
\ddot{x}_{2}(t) + \Delta_{2}\dot{x}_{2}(t) + \omega_{2}^{2}x_{2}(t) + \sqrt{M_{1}M_{2}^{-1}} \ \gamma \ \dot{x}_{1}(t) = F_{2}(t), \\
\gamma = G_{c}(\sqrt{M_{1}M_{2}})^{-1}
\end{cases}$

T.D. Scharton, R.H. Lyon, Power flow and energy sharing in random vibration. *Journal of the Acoustical Society of America*, 43, 1332-1343, 1968.

• Time-averaged power exchanged by two oscillators Case of white noise forces in $[\Omega_1, \Omega_2]$

Oscillators excited by white noise force in the frequency band $[\Omega_1, \Omega_2]$

$$P_{12} \approx \beta \left(E_1 - E_2 \right)$$

if
$$\begin{cases} \omega_1 \in [\Omega_1, \Omega_2] \\ \omega_2 \in [\Omega_1, \Omega_2] \end{cases}$$

Oscillator energy ratio E2/E1 versus the natural frequency of oscillator 2. Natural frequency of oscillator 1 = 1000 Hz.

Expression adapted for **resonant** modes (only)

• Energy sharing between mode *p* and mode *q*

• Modal energy equations of motions (for the two coupled subsystems):

$$\begin{cases} \Pi_{inj}^{1p} = \left(\omega_p^1 \eta_p^1 + \sum_{q'=1}^{N_2} \beta_{pq'}^{12}\right) E_p^1 & -\sum_{q'=1}^{N_2} \beta_{pq'}^{12} E_{q'}^2, \quad \forall p \in [1, \dots, N_1], \\ \\ \Pi_{inj}^{2q} = & -\sum_{p'=1}^{N_1} \beta_{p'q}^{12} E_{p'}^1 + \left(\omega_q^2 \eta_q^2 + \sum_{p'=1}^{N_1} \beta_{p'q}^{12}\right) E_q^2, \quad \forall q \in [1, \dots, N_2]. \end{cases}$$

→ $N_1 + N_2$ equations, $N_1 + N_2$ unknowns → SmEdA

• Subsystem energies:

$$E_1 = \sum_{p=1}^{N_1} E_p^1, \quad E_2 = \sum_{q=1}^{N_2} E_q^2$$

• Response in term of physical quantities:

$$< V_{\alpha}^{2} >= \frac{E_{\alpha}}{M_{\alpha}} < p_{\alpha}^{2} >= \frac{\rho_{0} c_{0}^{2} E_{\alpha}}{V}$$

Relation with SEA-like (Statistical Energy Analysis)
 Modal energy equipartition assumption:

$$E_p^1 = e_1, \forall p \in [1, ..., N_1], E_q^2 = e_1, \forall q \in [1, ..., N_2]$$

→ SEA equations (by summing the SmEdA equations):

$$\begin{cases} \Pi_{inj}^{1} = \omega_{c} \eta_{1} E_{1} + \omega_{c} \left(\eta_{12} E_{1} - \eta_{21} E_{2} \right), \\ \Pi_{inj}^{2} = \omega_{c} \eta_{2} E_{2} + \omega_{c} \left(\eta_{21} E_{2} - \eta_{12} E_{1} \right), \end{cases}$$

with the subsystem injected powers,

$$\Pi^{1}_{inj} = \sum_{p=1}^{N_{1}} \Pi^{1p}_{inj}, \ \Pi^{2}_{inj} = \sum_{q=1}^{N_{2}} \Pi^{2q}_{inj}, \ \text{and},$$

the SEA-like coupling loss factor (CLF),

$$\eta_{12} = \frac{\sum_{p=1}^{N_1} \sum_{q=1}^{N_2} \beta_{pq}^{12}}{N_1 \omega_c}$$

depanding only on the modal information of each uncoupled subsystem.

(1) Subsystems with low modal overlap

- →Low damping
- →Low modal density
- →Mid-frequency domain

Example on a case studied in the literature : Two coupled beams with varying damping FF. YAP, J. WOODHOUSE - Investigation of damping effects on SEA of coupled structures, JSV, 197 (1996)

L. Maxit, J.L. Guyader - Extension of SEA model to subsystems with non uniform modal energy distribution. *Journal of Sound and Vibration*, 265 (2003) 337-358.

(2) Heterogeneous subsystems

→ Even if a impedance mismatch has not been considered in the SmEdA substructuring, it is taken into account in the model (through the spatial mode shapes).

L. Maxit, J.L. Guyader - Extension of SEA model to subsystems with non uniform modal energy distribution. *Journal of Sound and Vibration*, 265 (2003) 337-358.

(3) Spatially localised excitations

Two coupled plates excited by a point force

Plate energy ratio (dB) for 12 different excitation points. 1000 Hz third octave band results.

Model energy distribution of the excited plate. Case of the excitation order 7.

Model energy distribution of the receiving plate. Case of the excitation order 7.

L. Maxit – Reformulation and extension of SEA model by taking the modal energy distribution into account, PhD thesis, INSA-Lyon, 2000 (in French).

(4) Hybrid SEA/SmEdA model

Mo: Modal overlap (= Damping bandwidth x modal density).

"Rain on the roof" excitation on plate 1.

Hybrid SEA-SmEdA model:

- plate 1 and 4 described by SEA
- plate 2 and 3 described by SmEdA

Energy ratio E4/E1 (dB) for each third octave band

L. Maxit – Reformulation and extension of SEA model by taking the modal energy distribution into account, PhD thesis, INSA-Lyon, 2000 (in French).

L. Maxit, J.L. Guyader - Estimation of sea coupling loss factors using a dual formulation and fem modal information, part 2 : numerical applications. *Journal of Sound and Vibration*, 239 (2001) 931-948.

Comparison with virtual experiments (FEM simulation)

Vibration transmission through
car firewall (RENAULT - 2001)Vibration transmission through
car floor

Sound radiation from car structure (2009)

Vibration transmission through car floor (FIAT - 2002)

N. Totaro, C. Dodard, J.L. Guyader, SEA Coupling Loss Factors of Complex Vibro-Acoustic Systems, JVA, ASME, 2009 ³¹

Industrial applications developped in the past

(6) Estimation of the local response (spatial energy distribution)

Modal expansion for a given structural subsystem:

$$v(M,\omega) = \sum_{p=1}^{N_i} j\omega a_p(\omega) \widetilde{W}_p(M)$$

Time-averaged square velocity at point M:

$$< v^{2}(M, \omega) >_{\Delta \omega} = \sum_{p=1}^{N_{i}} < \omega^{2} |a_{p}(\omega)|^{2} >_{\Delta \omega} \left[\widetilde{W}_{p}(M) \right]^{2} + \sum_{p=1}^{N_{i}} \sum_{\substack{q=1 \ \neq r}}^{N_{i}} < \omega^{4} a_{p}(\omega) = \sum_{\Delta \omega} \widetilde{W}_{p}(M) \widetilde{W}_{q}(M)$$
Neglected

Modal energy in function of the modal amplitude:

$$E_{p} \approx 2 < E_{p, Kinetic} >_{\Delta \omega} = <\omega^{2} |a_{p}(\omega)|^{2} >_{\Delta \omega} M_{p}$$

Time-averaged square velocity at point M in function of the modal energies:

$$< v^2(M, \omega) >_{\Delta \omega} \approx \sum_{p=1}^{N_i} \frac{E_p}{M_p} [\widetilde{W}_p(M)]^2$$

... similar for an acoustic subsystem.

N. Totaro, J.L. Guyader - Extension of the statistical modal energy distribution analysis for estimating energy density in coupled subsystems. Journal of Sound and Vibration, 331 (2012) 3114-3129.

Illustration of SmEdA application: Estimation of the pressure at the driver ear of a truck cab when the floor is excited by a point force

FUI-FEDER project (2012-2015)

Experimental set-up: supported truck cab excited by a mechanical force on a girder

- → dominant path (assumption): floor radiating into the cavity
- → First studied configuration: structure body in white

Y. Gerges, et al., Mid-frequency vibroacoustic modeling of an innovative lightweight cab – floor/cavity Interaction, proceeding of VISHNO, Aix en provence, France, June 2014.

FE mesh of the truck cavity

FE mesh of the floor

Y. Gerges, H.D. Hwang, et al., Vibroacoustic modeling of a trimmed truck cab in the mid frequency range. Proceeding of Internoise 2015, San Francisco, USA, August 2015

Example of intermodal coupling factors between the floor and the cavity (400 Hz third octave band):

SmEdA process for local energy estimation:

- Normal mode analysis on the uncoupled-subsystems (i.e. floor, cavity)
- Intermodal coupling factors calculations
- Modal energies estimation (by resolving the SmEdA equations)
- Local pressure estimation from the modal energies

Y. Gerges, H.D. Hwang, et al., Vibroacoustic modeling of a trimmed truck cab in the mid frequency range. Proceeding of Internoise 2015, San Francisco, USA, August 2015

SmEdA overview

- An extension of SEA by relaxing the modal energy equipartition assumption;
 - → application in the mid-frequency range (low modal overlap subsystems, local excitation, heterogeneous subsystems)
 - ➔ analysis of the modal energy transfers
 - → hybrid SEA/SmEdA models
- A formulation based on the knowledge of the subsystem modes;
 - → Used of subsystem Finite Element Models (FEM) for complex geometry
 - ➔ Useful tool for studying the different assumptions of the SEA (or SmEdA) fundamentals (see chapter 4, [*])
- A basic deterministic formulation
 - → Extension for including non-parametric uncertainties (like SEA)
 - Framework for future developments to include different degrees of uncertainties

IV. Extension of SmEdA to non-resonant transmission

Why there is an issue for the energy transmission between the non-resonant modes (in the SEA and SmEdA models)?

Oscillators excited by white noise force in the frequency band $[0,\infty[$

Oscillators excited by white noise force in the frequency band $[\Omega_1, \Omega_2]$

$$P_{12} \approx \beta \left(E_1 - E_2 \right)$$

if
$$\begin{cases} \omega_1 \in [\Omega_1, \Omega_2] \\ \omega_2 \in [\Omega_1, \Omega_2] \end{cases}$$

Oscillator energy ratio E2/E1 versus the natural frequency of oscillator 2. Natural frequency of oscillator 1 = 1000 Hz.

IV. Extension of SmEdA to non-resonant transmission

Context: Evaluation and analysis of the noise transmission through truck cab structures

<u>Question 1:</u> What is the TL of the floor when the emitting cavity is the engine compartment and the receiving cavity is the truck cabin ?

→ Influence of the sizes and shapes of the cavities on the noise transmission

→ Interest for a predictive method to estimate the TL of complex structures taking the behavior of small cavities and the structure geometry into account.

Estimation of the TL using Statistical Energy Analysis (SEA) model (for high frequency).

- → For f>fc (i.e.resonant transmission), the classical SEA model can be applied.
- ➔ For f<fc, Crocker and Price, JSV 9 (1969) proposed a modified SEA model to take the non resonant transmission into account.

<u>Question 2:</u> How can we estimated this CLF for a structure with a complex geometry?

IV. Extension of SmEdA to non-resonant transmission

(Black, resonant transmission, red, non resonant transmission)

The Dual Modal Formulation (DMF) allows us to write the matrix system:

$$\begin{bmatrix} Z_{11} & -j\omega W_{12} & 0 \\ +j\omega W_{12}^{*} & Z_{22} & +j\omega W_{23}^{*} \\ 0 & -j\omega W_{23} & Z_{33} \end{bmatrix} \begin{bmatrix} \Gamma_1 \\ \Gamma_2 \\ \Gamma_3 \end{bmatrix} = \begin{bmatrix} Q_1 \\ 0 \\ 0 \end{bmatrix}$$

Considering two sets of modes for the structure: the resonant (R) and the non-resonant (NR) gives:

Z_{11}	$-j\omega W_{12}^{NR}$	$-j\omega W_{12}^R$	0	$\left[\Gamma_1 \right]$	$\left[Q_{1} \right]$
$+j\omega W_{12}^{NR*}$	$Z_{22}^{\scriptscriptstyle NR}$	0	$+j\omega W_{23}^{NR*}$	Γ_2^{NR}	0
$+j\omega W_{12}^{R^*}$	0	Z_{22}^R	$+j\omega W_{23}^{R^*}$	Γ_2^R	0
0	$-j\omega W_{23}^{NR}$	$-j\omega W_{23}^R$	Z ₃₃	$\left[\Gamma_3 \right]$	

Achieving a condensation on the NR modes and assuming mass controlled behaviour for these modes, we obtain:

$$\begin{bmatrix} Z_{11} & -j\omega W_{12}^{R} & -W_{12}^{NR} W_{23}^{NR*} \\ +j\omega W_{12}^{R*} & Z_{22} & +j\omega W_{23}^{R*} \\ -W_{23}^{NR} W_{12}^{NR*} & -j\omega W_{23}^{R} & Z_{33} \end{bmatrix} \begin{bmatrix} \Gamma_{1} \\ \Gamma_{2} \\ \Gamma_{3} \end{bmatrix} = \begin{bmatrix} Q_{1} \\ 0 \\ 0 \end{bmatrix}$$

L. Maxit, K. Ege, N. Totaro, J.L. Guyader - Non resonant transmission modelling with SmEdA for evaluating the transmission loss of complex structures, *Journal of Sound and Vibration, 333 (2014) 499-519.*

Direct intermodal coupling factor between the two cavities:

$$\beta_{pr} = \left(\sum_{q \in \mathcal{Q}^{NR}} W_{pq} W_{rq}\right)^{2} \left\{ \frac{\left(\omega_{p} \eta_{p} + \omega_{r} \eta_{r}\right)}{\left[\left(\omega_{p}\right)^{2} - \left(\omega_{r}\right)^{2}\right]^{2} + \left(\omega_{p} \eta_{p} + \omega_{r} \eta_{r}\right) \left[\omega_{p} \eta_{p} \left(\omega_{r}\right)^{2} + \omega_{r} \eta_{r} \left(\omega_{p}\right)^{2}\right] \right\}}$$

- p: modes of cavity 1
- r : modes of cavity 2
- q : non-resonant modes of the plate
- W_{pq} W_{rq} are the intermodal works between modes of the cavities and non-resonant modes of the structure

SEA coupling loss factor between the 2 cavities:

$$\eta_{C1-C2} = \frac{\sum_{p \in P} \sum_{q \in R} \beta_{pr}}{P \omega_c}$$

IV. Extension of SmEdA to non-resonant transmission

Acoustic transmission between two "Small" cavities

	400 Hz	500 Hz	630 Hz	800 Hz	1 kHz	1.25kHz	1.6 kHz	2 kHz	2.5 kHz
Р	5	6	12	22	41	71	149	263	535
$Q^{\scriptscriptstyle NR}$	46	59	75	96	124	157	198	251	322
$Q^{\scriptscriptstyle R}$	13	16	21	28	33	42	52	70	83
R	4	4	12	21	32	67	129	231	472
	3.15kHz	4 kHz	5 kHz	6.3 kHz	8kHz	10kHz	12.5kHz	16kHz	20kHz
Р	1033	1998	3982	7815	15490	30672	60818	121228	236518
$Q^{\scriptscriptstyle NR}$	406	515	655	829	1049	1331	1682	2122	2687
$Q^{\scriptscriptstyle R}$	108	139	173	219	281	350	439	564	703
R	909	1766	3483	6859	13560	26876	53361	106085	209151

Subsystem mode numbers for each third octave band

IV. Extension of SmEdA to non-resonant transmission

Intermodal works (Third octave band 1000 Hz)

1mm-thick steel plate (critical frequency ~ 11.7 kHz)

Energy Noise Reduction versus third octave band:

SmEdA results for different plate DLF: cross 10%; circles, 1% ; square, 0.1%. Comparison of four calculations:

solid line, reference;

circles, SmEdA taking the NR plate modes into accour dashed line, DMF without NR plate modes;

diamonds, SmEdA without NR plate modes.

IV. Extension of SmEdA to non-resonant transmission

x, SmEdA; o, Simplified SmEdA

Validation of the SmEdA process with FEM modes

FEM mesh size criterion: flexural wavelength / 6

→ 5 % difference between the analytical and FEM modal frequencies

IV. Extension of SmEdA to non-resonant transmission

Example of results for stiffened plate: Plate thickness: 1mm Rib cross-section: square (5mm x 5mm) Rib spacing: 50 mm

V. Methodology for including the effect of dissipative treatments in SmEdA

Vibro-acoustic modelling of the truck cab including **dissipative treatments**

- Viscoelastic layers (Damping layer)

- Acoustic absorbing materials (trim, foarm)

→ How to take the dissipative effect of these materials into account with a SmEdA model?

Modal energy equations of motion for two subsystems (SmEdA):

$$\begin{cases} \Pi_{inj}^{1p} = \left(\omega_{pq}^{1} + \sum_{q'=1}^{N_{2}} \beta_{pq'}^{12}\right) E_{p}^{1} & -\sum_{q'=1}^{N_{2}} \beta_{pq'}^{12} E_{q'}^{2}, \quad \forall p \in [1, ..., N_{1}], \\ \Pi_{inj}^{2q} = & -\sum_{p'=1}^{N_{1}} \beta_{p'q}^{12} E_{p'}^{1} + \left(\omega_{q}^{1} \eta_{q}^{2} + \sum_{p'=1}^{N_{1}} \beta_{p'q}^{12}\right) E_{q}^{2}, \quad \forall q \in [1, ..., N_{2}]. \end{cases}$$

with the intermodal coupling factors:

$$\beta_{pq}^{12} = \frac{W_{pq}^{12}}{M_{p}^{1} (\omega_{q}^{2})^{2} M_{q}^{2}} \left\{ \frac{\omega_{p} \eta_{p}^{1} (\omega_{q}^{2})^{2} + \omega_{p}^{2} \eta_{q}^{2} (\omega_{p}^{1})^{2}}{\left[(\omega_{p}^{1})^{2} - (\omega_{q}^{2})^{2} \right]^{2} + (\omega_{p} \eta_{p}^{1}) + \omega_{p}^{2} \eta_{q}^{2} (\omega_{p} \eta_{p}^{1}) (\omega_{q}^{2})^{2} + \omega_{p}^{2} \eta_{q}^{2} (\omega_{p}^{1})^{2} \right] \right\}$$

Dissipative effect \rightarrow Modal damping loss factors \rightarrow $\eta_p^1, \forall p \in [1,...,N_1]$ $\eta_q^2, \forall q \in [1,...,N_2]$ Dissipative powers Intermodal coupling factors

Illustration of the methodology on the validation test cases PhD thesis H.D. HWANG (2011-2015)

Experimental set-up of the validation case (Rectangular "clamped" plate radiating into a rectangular cavity with "rigid" wall)

H.D. Hwang et al.- A methodology for including the effect of a damping treatment in the mid-frequency domain using SmEdA method, *Proceedings of 20th International Congress on Sound and Vibration*, Bangkok, Thailand, July 2013. 8p.

Illustration of the methodology on the validation test case

(MOVISAND software)

1st step: Characterisation of the equivalent dissipative material

Equivalent Fluide ($\rho_{eq}(f), c_{eq}(f), \eta_{eq}(f)$) (Acoustic tube with 2 cavities method)

Equivalent Single Layer Model for Viscoelastic Materials

One root associated to the dominant transverse motion \rightarrow bending stiffness of the equivalent layer

$$E_{eq}^{*} = \frac{12(1 - v_{eq}^{2})B}{h_{1}^{3}}$$
• Reduced FEM computation compared to 3D model
• Method implemented in *MOVISAND* software

Equivalent Fluid for Porous Materials

Biot's model (Theory of poroelasticity)

- Solid (u^s) + fluid (u^f)
- One shear, two compression waves
- Macroscopic properties (φ, σ, α, Λ, Λ)

Equivalent fluid model

- Rigid solid assumption (u^s = 0)
- Material fixed to a rigid surface
- Characterized as a fluid ($c_{eq'} \rho_{eq}$)
- Kundt tube measurement $\rightarrow Z_c$ & k_{eq}

Acoustic tube measurement with the two cavities method, H. Utsuno, et al. *JASA*,86(2), 1989

thermal dissipative effect

Illustration of the methodology on the validation test case

!!! Equivalent material properties depend on the third octave band **!!!**

Illustration of the methodology on the validation test case

<u>3nd step:</u> Normal modes extraction and evaluation of the modal damping loss factors (from the imaginary part of the FE matrices) NASTRAN DMAP

<u>4th step:</u> Calculation of the modal coupling loss factors (MCLFs)

<u>5th step:</u> Calculation of the modal energies and the subsystem energy from SmEdA equations

 $(\omega_q, M_q, \mathbf{p}_q)$ $(\omega_p, M_p, \mathbf{W}_p)$ **MSE** method **MSKE** method $\frac{\mathrm{Im}\{\mathrm{W}_{p}\overline{\mathrm{K}}\mathrm{W}_{p}\}}{\mathrm{Re}\{\mathrm{W}_{n}\overline{\mathrm{K}}\mathrm{W}_{n}\}}$ η_p (see [PhD Hwang]) $\beta_{pq} = f(\omega_p, M_p, \mathbf{W}_p, \eta_p, \omega_a, M_a, \mathbf{p}_a, \eta_a)$ E_p, E_a E_{t1}, E_{t2}

H.D. Hwang, Extension of the SmEdA method by taking into account dissipative materials in medium frequency, PhD thesis, INSA Lyon, France, 2015.

Four test cases:

- → Comparison of the no-treatment case to the treatment cases
- → To study the influence of the treatments on the energy transmission

V. Methodology for including the effect of dissipative treatments in SmEdA

provided by ACOEM (Dynamic Material Analyser)

V. Methodology for including the effect of dissipative treatments in SmEdA

<u>2nd STEP</u> FE modeling of the subsystems with equivalent parameters

•4,031,412 solid elements•6 kHz band•Rigid walls

Experimental validation: High-resolution modal analysis based on ESPRIT algorithm

III For heavily damped modes,
 SNR is problematic experimentally.
 →ESPRIT effective for lightly damped modes!!!

K. Ege, et al., High resolution modal analysis, Journal of Sound and Vibration, 325 (2009) 852-869

V. Methodology for including the effect of dissipative treatments in SmEdA

Computation of the Modal coupling factors

$$\beta_{pq}^{12} = \frac{W_{pq}^{12}}{M_{p}^{1} (\omega_{q}^{2})^{2} M_{q}^{2}} \left\{ \frac{\omega_{p}^{1} \eta_{p}^{1} (\omega_{q}^{2})^{2} + \omega_{q}^{2} \eta_{q}^{2} (\omega_{p}^{1})^{2}}{\left[(\omega_{p}^{1})^{2} - (\omega_{q}^{2})^{2} \right]^{2} + (\omega_{p}^{1} \eta_{p}^{1} + \omega_{q}^{2} \eta_{q}^{2}) \left[\omega_{p}^{1} \eta_{p}^{1} (\omega_{q}^{2})^{2} + \omega_{q}^{2} \eta_{q}^{2} (\omega_{p}^{1})^{2} \right] \right\}$$

Modal coupling (β_{pq}^{12}) at 1 kHz

V. Methodology for including the effect of dissipative treatments in SmEdA

<u>5th STEP</u> Computation of the subsystem energies

Experimental Validation: Subsystem energy ratio (E_2/E_1)

At low frequency, discrepancies due to the boundary condition? Next future: Application to the truck cab with trims (Y. Gerges) Recent works achieved by ACOEM (acoustic consulting company): - Automation of the numerical process through an in-house code (developed under

ANSYS-APDL environment)

- Benchmark on a mock-up of a nuclear power plant structure

- Applications on industrial buildings

P. VOUAGNER, L. MAXIT, C. THIRARD, C. DESLOT, J.L. GUYADER, Modélisation de la propogation du bruit solidien dans les structures industrielles. CFA 2014, Poitier, France, April 2014.

SmEdA overview

- SEA energy equipartition assumption relaxed
 - \rightarrow Extension to low modal overlap subsystems
 - \rightarrow Extension to non-homogeneous subsystems (with local energy description)
- Method based on the uncoupled subsystems modes;
 - → Description of subsystems with complex geometry/mechanical properties (using FEM) → Description of dissipative treatments
- Non-resonant transmission modelling
 - → Prediction of TL of complex structures in mid-frequency
- Hybrid SEA/SmEdA model

Future research on SmEdA

- Uncertainty (screening technique, propagation of uncertainties, variance estimation,...)
- Tools for energy transfer analysis (using the graph theory: tomorrow, Oriol Guasch,...)

Statistical modal Energy distribution Analysis

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Y. Gerges, H.D. Hwang, et al., Vibroacoustic modeling of a trimmed truck cab in the mid frequency range. Proceeding of Internoise 2015, San Francisco, USA, August 2015

Thank you for your attention